

Emergence of spacetime from string-theoretic Matrix Models

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Motivation



↪ Inflation: An early phase of **accelerated expansion**.

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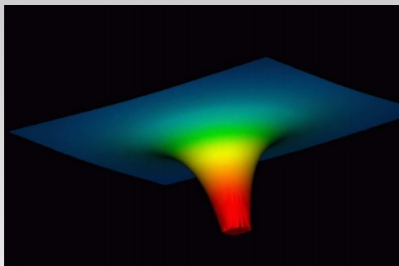


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Credit: Pablo Laguna

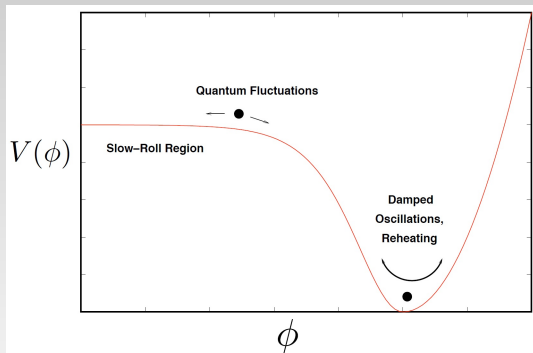
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Alternate description of the early-universe from **Matrix Theory**.

Successful scenarios of early-universe cosmology



↔ SBB cosmology **successfully** explains:

- ✓ Expansion of the Universe: **Hubble Law**
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- ✓ Abundance of light elements: **BBN**

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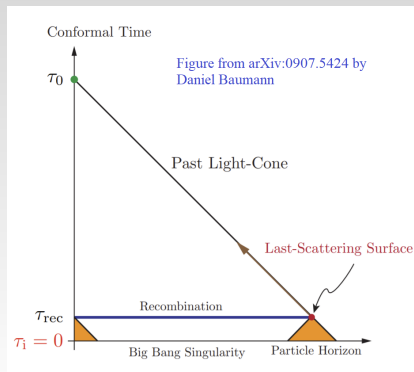
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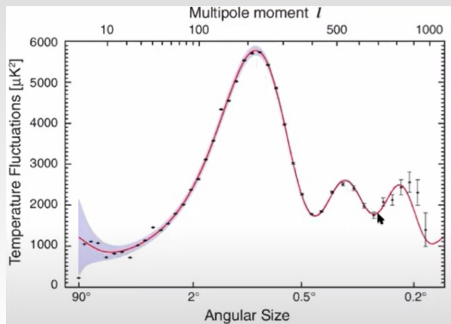
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- Solution to the standard **horizon problem** \Leftarrow Hubble radius **smaller** than causal horizon is all that is needed. [Brandenberger, 2011]
- Solution to the standard **flatness problem** \Leftarrow Spatially flat universe.
- Origin of structure \Leftarrow **Scale invariant** power spectrum of adiabatic perturbations. [Sunyaev & Zel'Dovich; Peebles & Yu, 1970]
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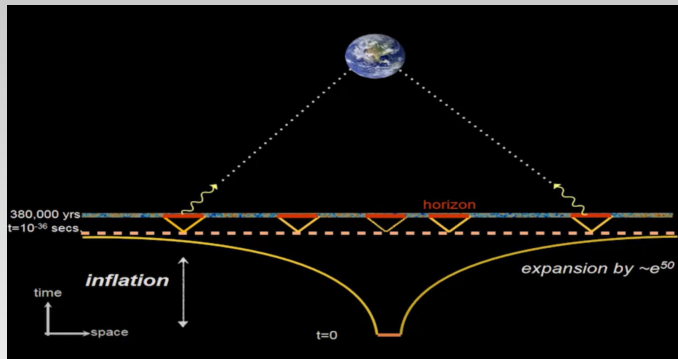


Photo credit: P. Adshead

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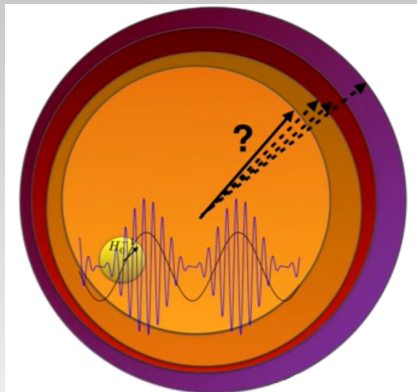
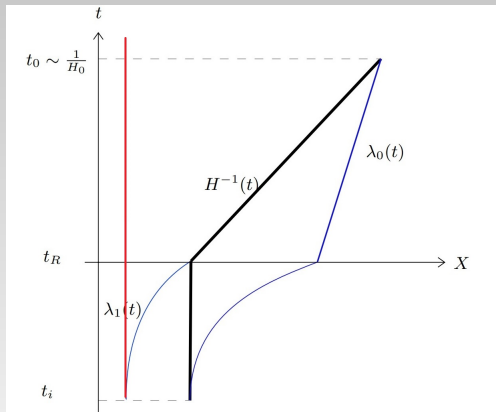


Photo credit: S. Shandera

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[Bedroya, Brandenberger, LoVerde, Vafa, 2019]

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Symmetries of QG, *e.g.*, **String Dualities**, can play a pivotal role in this!
New symmetries (T-duality) \Leftrightarrow New states (**Winding modes**)

Alternatives

↪ Alternate descriptions of early-universe cosmology:

- Examples: *String Gas Cosmology* [Brandenberger & Vafa, 1989], *Ekpyrotic bounce* [Khoury, Ovrut, Steinhardt & Turok, 2001], *Early phase of topological gravity* [Agrawal, Gukov, Obied & Vafa, 2020], *CPT-symmetric universe* [Boyle & Turok, 2018-22], *Matter bounce* [Brandenberger, Wands, Wilson-Ewing, Cai & Wilson-Ewing, ...], , ...

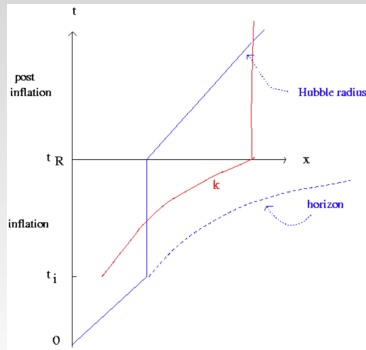
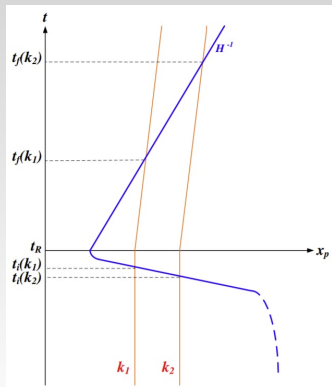


Figure Credit: R. Brandenberger



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where $D_t \equiv \partial_t - i[A_t, \cdot]$ and $N \times N$ bosonic matrices $A(t), X_i(t)$ ($i = 1, \dots, d$) and $\psi_\alpha(t)$ ($\alpha = 1, \dots, p$).

✓ **Quantum Mechanical model** of 9 bosonic $U(N)$ matrices and their 16 fermionic superpartners.

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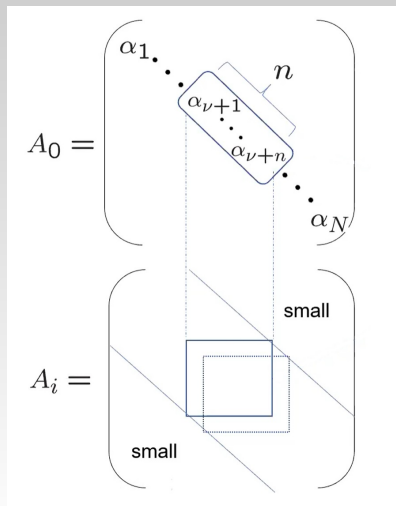
↔ A **consistent, principled** *top-down* approach to early-universe cosmology with the promise of **rich phenomenology**.

- **No Fock-space EFT** description: **No cosmological constant problem!**
- **UV-complete:** Eigenvalues never become **trivial** ⇒ **No singularities!**

Emergent (3 + 1)-d spacetime: Numerical evidence

[Aoki, Hirasawa, Ito, Kim, Nishimura, Tsuchiya, ...]

↪ Lorentzian **IKKT** model: $Z \sim \int dA d\Psi e^{iS_{\text{IKKT}}}$



✓ **Diagonalize** $A_0 : \alpha_1 < \dots < \alpha_N$.

✓ Define **time** via *coarse-graining*:

$$t(\nu) := \frac{1}{n} \sum_{i=1}^n \alpha_{\nu+i}, \quad \nu = 1, \dots, N - n$$

✓ *Non-trivial* to obtain dynamical **band-diagonal structure!**

→ **Time-dep** $n \times n$ spatial matrices:

$$(\bar{A}_i)_{I,J}(t(\nu)) := (A_i)_{\nu+I, \nu+J}$$

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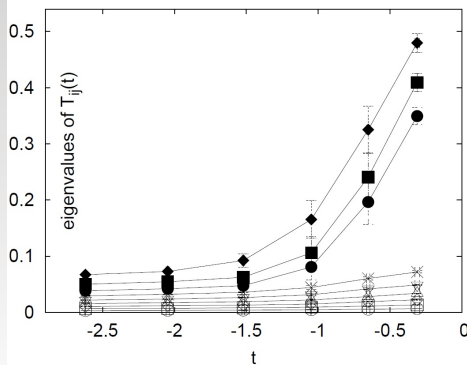


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↪ *Emergence* of 3 large spatial dimensions:



✓ As order parameter, define **moment of inertia tensor**

$$T_{ij}(t) := \left\langle \frac{1}{n} \text{Tr} \bar{A}_i(t) \bar{A}_j(t) \right\rangle$$

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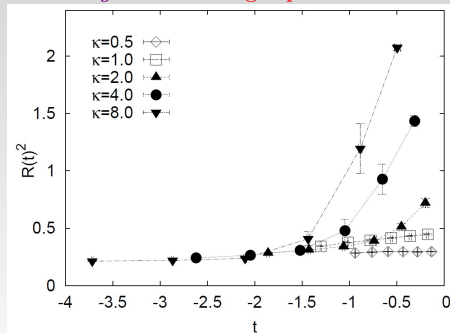


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✓ Extent of a given spatial dimension parameter:

$$x_i^2(t) := \left\langle \frac{1}{n} \text{Tr} \bar{A}_i(t)^2 \right\rangle$$

→ Total extent of space parameter:

$$R^2(t) = \sum_{i=1}^9 x_i^2(t)$$

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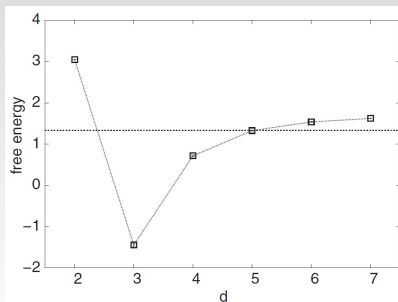
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Takeaway: A $(3 + 1)$ -d universe **emerges dynamically** from non-perturbative (matrix model) description of superstring theory

Coarse-graining spacetime from abstract matrices

[S.B., Brandenberger & Laliberté, 2206.12468 (JHEP)]

↔ Emergent time identified from the diagonal elements of A_0 matrix, ordered as $A_0 = \text{diag}(t_1, \dots, t_N)$ with $t_i > t_j$ for $i > j$.

↔ In the diagonal A_0 basis, the eigenvalues of the A_i matrices **decay** when moving away from the diagonal ⇒ **Band-diagonal structure**.

✓ $t_{\max} \sim \sqrt{N}$, and discrete time eigenvalues scale as: $\Delta t \sim 1/\sqrt{N}$.

✓ In the $N \rightarrow \infty$ limit, *emergent* continuous and infinite time.

✓ Total **physical extent of space**: $l_{\text{phys}} \sim \sqrt{N}$.

✓ Continuous and infinite extent of space when $N \rightarrow \infty$!

↔ Emergent metric: $g_{ij}(n, t) = \mathcal{A}(t) \delta_{ij}$, Assumption: $SO(3)$ symmetry.

No Flatness Problem! Independent of isotropy assumption.

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↔ Emergent metric: $g_{ij}(n, t) = \mathcal{A}(t) \delta_{ij}$, **Assumption**: $SO(3)$ symmetry.

No Flatness Problem! Independent of isotropy assumption.

Coarse-graining spacetime from abstract matrices

[S.B., Brandenberger & Laliberté, 2206.12468 (JHEP)]

↔ Emergent time identified from the diagonal elements of A_0 matrix, ordered as $A_0 = \text{diag}(t_1, \dots, t_N)$ with $t_i > t_j$ for $i > j$.

↔ In the diagonal A_0 basis, the eigenvalues of the A_i matrices **decay** when moving away from the diagonal ⇒ **Band-diagonal structure**.

✓ $t_{\max} \sim \sqrt{N}$, and discrete time eigenvalues scale as: $\Delta t \sim 1/\sqrt{N}$.

✓ In the $N \rightarrow \infty$ limit, **emergent continuous and infinite time**.

✓ Total **physical extent of space**: $l_{\text{phys}} \sim \sqrt{N}$.

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Cosmology from a thermal state in BFSS

↔ Consider a **thermal state in the BFSS model**.

↔ Start with the Euclidean BFSS model and consider its **compactification on a thermal circle**: BFSS $\xrightarrow{T \rightarrow \infty}$ IKKT (**natural** to assume thermal state)

$$\rightsquigarrow S_{\text{BFSS}}(\beta) = \frac{1}{2g^2} \int_0^\beta dt \text{Tr} \left\{ (D_t X_i)^2 - \frac{1}{2} [X_i, X_j]^2 + \text{fermions} \right\}, \quad \beta = 1/T.$$

↪ Fourier expand fields as $X_i(t) = \sum_n X_i^n e^{in\omega t}$ in Matsubara frequencies $\omega = 2\pi T$.

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✓ One can just as well start with a thermal state in the IKKT model (by compactifying the A_0 matrix through the method of images) and use that to derive cosmology ($T \rightarrow 1/T$). We find the same background and a scale invariant spectrum, although with slight $\mathcal{O}(1)$ factors different in the amplitude. [Laliberté & S.B., 2304.10509 (JHEP)]

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↪ Calculate **thermodynamic quantities** in the *Euclidean BFSS model* and expand in the **high T limit** (dimensionless expansion parameter: $\sqrt{g^2 N / T^3}$)

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✓ Integrate over **non-zero modes alone** to arrive at above results since we want to expand **around the IKKT background**. After integrating out non-zero modes using perturbation theory, the **leftover integration over the zero modes** can be thought of as taking the expectation value of connected Green's functions using the bosonic part of the IKKT action.

✓ Use approximations to evaluate higher-order moments of zero-modes.

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$\chi_1 := \langle \frac{1}{N} \text{Tr} (A_i)^2 \rangle_{\text{IKKT}} \propto R^2$, $\chi_2 := \langle \frac{1}{N} \text{Tr} (F_{ij})^2 \rangle_{\text{IKKT}}$ evaluated in S_0 .

✓ Integrate over **non-zero modes alone** to arrive at above results since we want to expand **around the IKKT background**. After integrating out non-zero modes using perturbation theory, the **leftover integration over the zero modes** can be thought of as taking the **expectation value of connected Green's functions** using the bosonic part of the **IKKT action**.

✓ Use **approximations** to evaluate higher-order moments of zero-modes.

✓ C_V **positive** for $d = 3, p = 4$ and thermodynamics is **well-defined**.

→ We have all the quantities **required** C^{00}_{00} , C^{ij}_{ij} to calculate spectrum of **cosmological perturbations** in the **BFSS thermal state**. [Kawahara, Nishimura & Takeuchi, 2007; S.B., Brandenberger & Laliberté, 2107.11512]

Cosmological observables from the BFSS theory



We find scale-invariant spectrum for *IR modes* of observational interest

[S.B., Brandenberger & Laliberté, 2107.11512]

$$P_{\zeta} \sim (\ell_s M_{\text{Pl}})^{-4} \mathcal{O}(1), \quad P_h \sim \alpha (\ell_s M_{\text{Pl}})^{-4} \mathcal{O}(1). \quad \text{SG: } \mathcal{A} \sim (\ell_s M_{\text{Pl}})^{-4}$$

↔ UV-modes for density perturbations have a Poisson spectrum ($\propto k^2$); distinct from inflation but not of (direct) observable consequence. Tensor spectrum only has a scale-invariant part.

✓ Phenomenologically relevant predictions to appear for tensor-to-scalar ratio ($\alpha \sim$ off-diagonal), abundance of PBHs (UV spectrum), large primordial B-fields (thermal state) and non-Gaussianities (bispectra). Small tilt from next order corrections and due to the SSB phase transition.

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- ✓ Single $U(N)$ matrix QM is known to be **equivalent** to $2d$ String Theory.

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• The **collective field formalism**:

$$\phi_k = \operatorname{Tr} (e^{ikM}) = \sum_{i=1}^N e^{ik\lambda_i}$$

$$\phi(x) = \int \frac{dk}{2\pi} e^{-ikx} \operatorname{Tr} (e^{ikM}) = \operatorname{Tr} (\delta(x - M)) = \sum_{i=1}^N \delta(x - \lambda_i)$$

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- The **collective field formalism**: Description in terms of **density eigenvalues** with the constraint $\int dx \phi(x) = N$

- ✓ Change of variables $M \rightarrow \phi$ leads to $(1+1) - d$ field theory:

$$H_\phi = \int dx \left[\frac{1}{2} \partial_x \pi(x) \phi(x) \partial_x \pi(x) + \frac{\pi^2}{6} \phi^3(x) - (\mu_F - V(x)) \phi(x) \right]$$

The collective field has a natural interpretation as a field in a higher dimensional theory [Das & Jevicki]

Collective field formalism for multi-matrix models



↔ A direct transformation to collective field for more than one matrices can only be implemented numerically → **Impossible** to solve the ‘Schwinger-Dyson’ equations analytically.



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$$\phi_C = \text{Tr} \left(M_1^a M_2^b M_3^c \dots \right)$$

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$$S_{\text{eff}} = \int d\tau \left[\frac{1}{2l_s} \left(\sum_i^N \dot{\lambda}_i^2 + \sum_i^N \dot{\rho}_i^2 \right) \right] - \frac{1}{2} \sum_{i<j, i=1}^N \text{Tr} \log \left(1 - \frac{1}{2l_s^4} \int ds |s - \tau| (\lambda_i(s) - \lambda_j(s))^2 \right)$$

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To get a **time-local effective action**, we need to add a **mass term** $\text{Tr} (m^2 Y^2)$ to the action. Then, to leading order:

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Possible to get a $(2+1)$ -d collective field action! Starting point for **connection to String Theory** → *Geometry from Entanglement*

[with Brandenberger, Dasgupta, Lei & Pasiiecznik]



$$\begin{aligned}
 H = & \frac{l_s}{2} \int dx dy [\partial_x \pi(x, y) \phi(x, y) \partial_x \pi(x, y) + \partial_y \pi(x, y) \phi(x, y) \partial_y \pi(x, y)] \\
 & + \frac{l_s \pi^2}{6} \int dx dy \phi^3(x, y) + \frac{l_s}{8} \int dx dy \frac{(\partial_y \phi(x, y))^2}{\phi(x, y)} \\
 & - \frac{1}{16m l_s^4} \int dx dy dx' dy' \phi(x, y) (x - x')^2 \phi(x', y') + \frac{l_s m^2}{2} \int dx dy y^2 \phi(x, y) \\
 & - \lambda \left[\int dx dy \phi(x, y) - N \right]
 \end{aligned}$$

- Has the **right expression** when we turn off the second matrix (compactifying y -direction to zero)!
- Yet to understand physical interpretation of taking the **massless limit** of this model to preserve connection to String Theory!

Conclusions

- ✓ It has been notoriously difficult to find accelerating solutions in string theory → Inflation as a coherent state over warped Minkowski? dS *does* have such an interpretation!

[S.B., Dasgupta & Tatar, *JHEP*, 2020; S.B., Dasgupta, Guo & Kulinich, *JHEP*, 2024]

- ✓ M-theory consistency rules out large classes of bounce models.

[Bernardo, S.B., Dasgupta Mir & Tatar, *Phys. Rev. Lett.*, 2021]

★ A new path towards a UV-complete paradigm for the early universe:

- ✓ Numerical evidence for the emergence of 3 large spatial dimensions from full String Theory.
- ✓ Analytically extract a coarse-grained time, space and metric.
- ✓ Thermal fluctuations ⇒ Scale invariant primordial perturbations.
- ✓ Horizon problem, Flatness problem and formation of structure from first-principles in a fundamental quantum gravity theory.
- ✓ No vacuum energy problem → Transition from non-geometric emergent phase to radiation dominated era. No cosmological constant.
- ✓ Possible to work with the full IKKT model alone and consider a thermal state in it. [Laliberté & S.B., *JHEP*, 2023]

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Looking ahead



- ★ A new path towards a UV-complete paradigm for the early-universe.

Ambitious, promising but a lot remains to be done:

- Details of non-geometric phase → Connection to **non-commutative geometry** emergent in the UV? [H. Steinacker; A. Chaney & A. Stern; ...]
- **Physical explanation** of the SSB phase? Connection with String Gas? Interestingly, non-geometric phase has $p = 0$ (quasi-static phase).
- ↪ Strings as solitonic states in Matrix models → Annihilation of string loops into radiation. **Collective field formalism!**
- Derive background cosmology from **BFSS model** (combine MC simulations and analytical insights). *Derive Friedmann equations!*
- A **gauge-invariant** notion of entanglement for matrices using the **Collective Field** → Area-law for toy models. [Frenkel & Hartnoll; Das, Kaushal, Mandal, Liu, Trivedi; Hampapura, Harper & Lawrence, ...]
- Observable consequences: **NG**, **PBHs** from the **Poisson part** of the UV spectrum, **Primordial B fields**, ...
- ⋮